

APS INVITED TALK N.Y.
Hidden Hamiltonian Structure
in Infinite Dimensional Systems
of Plasma Physics

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OVERVIEW

- I. Introduction
 - II. Background
 - III. Vlasov - Noncanonical
 - IV. Vlasov - Canonical
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My Poster 6R13 Wed.

Rel. M-V	Bialynicki - Birula Hubbard Amend	} 1S2 & 3
2-D	Spencer & Kaufman	} 8T8
MV	Marsden & Weinstein	} 8T9

II. Background

A. ORDINARY DIFF. EQS.

$$\mathcal{L}(q, \dot{q}) \xrightarrow{\text{LEGENDRE}} H(p, q)$$

$$\dot{q}_k = [q_k, H] \quad \dot{p}_k = [p_k, H]$$


$$k = 1, 2, \dots, N$$

where

$$[f, g] = \sum_k^N \left(\frac{\partial f}{\partial q_k} \frac{\partial g}{\partial p_k} - \frac{\partial g}{\partial q_k} \frac{\partial f}{\partial p_k} \right)$$

f & g are functions
of p_k & q_k .

Let $z^i = \begin{cases} q_k & k = i = 1, 2, \dots, N \\ p_k & i = k + N = k + 1, \dots, 2N \end{cases}$



$$[f, g] = \frac{\partial f}{\partial z^i} J^{ij} \frac{\partial g}{\partial z^j}$$

$$(J^{ij}) = \begin{bmatrix} 0 & I_N \\ -I_N & 0 \end{bmatrix}$$



Transforms as a
contravariant Tensor

Hamilton's Eqs. become

$$\dot{z}^i = [z^i, H] = J^{ij} \frac{\partial H}{\partial z^j}$$

Commutator Properties

- * $[f, g]$ is bilinear
- * $[f, g] = -[g, f]$
- * $[f, [g, h]] + [g, [h, f]] + [h, [f, g]] = 0$

Canonical Transformations

Preserve the form of J^{ij}

General Coordinate Transformations

Do Not!

$$\bar{J}^{st} = \frac{\partial \bar{z}^s}{\partial z^i} J^{ij} \frac{\partial \bar{z}^t}{\partial z^j}$$

General Coordinate Transformations

Do preserve Bracket Properties.

Converse Outlook: \bar{J}^{ij} has

properties but not skew-

diagonal. Can one find

Canonical Coordinates?

Yes! ($\det J^{ij} \neq 0$) Darboux
(1882)

DEF. A system $\dot{x}^i = F^i(x)$
is Hamiltonian if it
can be written in the
form $\dot{x}^i = [x^i, H]$ when
[,] makes vector
space of phase functions
into a Lie Algebra.

Because locally canonical
variables can be found

B. Field Equations (e.g. P.D.E. or I.E.)

$$[F, G] = \sum_{k=1}^n \int \left(\frac{\delta F}{\delta \eta_k} \frac{\delta G}{\delta \eta_k} - \frac{\delta F}{\delta \eta_k} \frac{\delta G}{\delta \eta_k} \right) dz$$

F & G are Functionals e.g.

$$H(s, v) = \int \frac{\rho v^2}{2} dz$$

$$\left. \frac{dF(\eta + \epsilon w)}{d\epsilon} \right|_{\epsilon=0} = \int \frac{\delta F}{\delta \eta} w dz = \left\langle \frac{\delta F}{\delta \eta} \mid w \right\rangle$$

$$[F, G] = \left\langle \frac{\delta F}{\delta \eta^i} \mid O^{ij} \frac{\delta G}{\delta \eta^j} \right\rangle$$

canonical $(O^{ij}) = \begin{bmatrix} 0 & I_N \\ -I_N & 0 \end{bmatrix}$, in

General Anti-self-Adjoint wrt $\langle \mid \rangle$
+ Jacobi condition

Def. If can write
a system of Field
Equations in the
form

$$\dot{X}_t^i = [X_t^i, H]$$

where $[,]$ has

properties, then call
Hamiltonian.

EXAMPLE K-dV Equation

$$\frac{\partial u}{\partial t} = + u \frac{\partial u}{\partial x} + \beta \frac{\partial^3 u}{\partial x^3}$$

$$H = \int \left(\frac{u^3}{6} - \beta \frac{u_x^2}{2} \right) dx$$

$$\frac{\delta H}{\delta u} = \frac{u^2}{2} + \beta u_{xx}$$

Gardner Bracket

$$[f, g] = \int \frac{\delta f}{\delta u} \frac{d}{dx} \frac{\delta g}{\delta u} dx$$

Note :

$$\frac{\partial u}{\partial t} = [u, H] = \frac{d}{dx} \left(\frac{u^2}{2} + \beta u_{xx} \right)$$

CAN CANONIZE !

III. Vlasov - Noncanonical

A. Vlasov-Poisson ($B=0$)

$$\frac{\partial f_\alpha(\underline{z}, t)}{\partial t} = - \underline{w}^\alpha \cdot \frac{\partial f_\alpha}{\partial \underline{z}}$$

$$\underline{w}^\alpha; \underline{z} \in \mathbb{R}^6$$

$$\underline{w}^\alpha = \left(\underline{v}, \frac{e_\alpha}{m_\alpha} \frac{d}{d\underline{x}} \sum_\beta \int_{\mathbb{R}^6} \nabla(\underline{x}|\underline{x}') f_\beta(\underline{z}'|t') d\underline{z}' \right)$$

$$\left(\text{Note: } \frac{\partial}{\partial \underline{z}} \cdot \underline{w}^{(\alpha)} = 0 \right)$$

like 2-D turbulence

$$H(f) = \sum_\alpha \frac{m_\alpha}{2} \int v^2 f_\alpha d\underline{z}$$

$$- \frac{1}{2} \sum_{\alpha, \beta} e_\alpha e_\beta \int \nabla(\underline{x}|\underline{x}') f_\alpha f_\beta d\underline{z} d\underline{z}'$$

$$[F, G] = \sum_{\alpha} \int \frac{f_{\alpha}(z)}{m_{\alpha}} \left\{ \frac{\delta F}{\delta f_{\alpha}}, \frac{\delta G}{\delta f_{\alpha}} \right\} dz$$

Where

$$\{f, g\} = \frac{\partial f}{\partial x} \cdot \frac{\partial g}{\partial u} - \frac{\partial f}{\partial u} \cdot \frac{\partial g}{\partial x}$$

Not difficult to see

$$\frac{df_{\alpha}}{dt} = [f_{\alpha}, H] \quad \square$$

Another way

$$\frac{df_{\alpha}}{dt} = -u \cdot \frac{\partial f_{\alpha}}{\partial x} - \frac{e_{\alpha}}{m_{\alpha}} E \cdot \frac{\partial f_{\alpha}}{\partial u}$$

$$\frac{\partial E}{\partial t} = - \sum_{\alpha} e_{\alpha} \int u f_{\alpha} du$$

$$H(f, \underline{E}) = \sum_{\alpha} \int \frac{m_{\alpha}}{2} v^2 f_{\alpha} d^3z + \frac{1}{2} \int E^2 d\underline{x}$$

$$[F, G] = \sum_{\alpha} \int_{\mathbb{R}^6} \frac{f_{\alpha}}{m_{\alpha}} \left\{ \frac{\delta F}{\delta f_{\alpha}}, \frac{\delta G}{\delta f_{\alpha}} \right\} d^3z$$

$$- \frac{e_{\alpha}}{m_{\alpha}} \int \left(\frac{\delta F}{\delta f_{\alpha}} \frac{\partial f_{\alpha}}{\partial v} \cdot \frac{\delta G}{\delta E} - \frac{\delta G}{\delta f_{\alpha}} \frac{\partial f_{\alpha}}{\partial v} \cdot \frac{\delta F}{\delta E} \right) d^3z$$

B. Maxwell-Vlasov

Maxwell's Eqs. in Vacuum

$$H = \int \frac{E^2 + B^2}{2} d\underline{x}$$

$$[F, G]_{\text{VM}} = \int \left(\frac{\delta F}{\delta E} \cdot \nabla \times \frac{\delta G}{\delta B} - \frac{\delta G}{\delta B} \cdot \nabla \times \frac{\delta F}{\delta E} \right) d\underline{x}$$

Jacobi Automatic

Whole Thing

$$\begin{aligned} [F, G] = & \sum_x \int \frac{f_x}{m_x} \left\{ \frac{\delta F}{\delta f_x}, \frac{\delta G}{\delta f_x} \right\} dx \\ & + \frac{e_x}{m_x} \int \left(\frac{\delta G}{\delta f_x} \frac{df_x}{dt} \frac{\delta F}{\delta E} - \frac{\delta F}{\delta f_x} \frac{df_x}{dt} \frac{\delta G}{\delta E} \right) dx \\ & + \frac{e_x}{m_x} \int \mathbf{B} \cdot \left(\frac{\partial}{\partial \mathbf{u}} \frac{\delta F}{\delta \mathbf{f}_x} \times \frac{\partial}{\partial \mathbf{u}} \frac{\delta G}{\delta \mathbf{f}_x} \right) dz \\ & + \int \left(\frac{\delta F}{\delta \mathbf{E}} \cdot \nabla \times \frac{\delta G}{\delta \mathbf{B}} - \frac{\delta G}{\delta \mathbf{B}} \cdot \nabla \times \frac{\delta F}{\delta \mathbf{E}} \right) dx \end{aligned}$$

Can prove Jacobi
Directly!!



$$\frac{\partial \underline{B}}{\partial t} = [\underline{B}, H] = -\nabla \times \underline{E}$$

$$\frac{\partial \underline{E}}{\partial t} = [E, H] = \nabla \times \underline{B} - \sum_{\alpha} e_{\alpha} \int \underline{v} f_{\alpha} d\underline{v}$$

$$\frac{\partial f_{\alpha}}{\partial t} = -\underline{v} \cdot \frac{\partial f_{\alpha}}{\partial \underline{x}} - \frac{e_{\alpha}}{m_{\alpha}} (\underline{E} + \underline{v} \times \underline{B}) \cdot \frac{\partial f_{\alpha}}{\partial \underline{v}}$$

Note: $\nabla \cdot \underline{B} = 0$

& $\nabla \cdot \underline{E} = \rho$

are initial conditions

$$[\nabla \cdot \underline{B}, H] = [\nabla \cdot \underline{E} - \rho, H] = 0$$

IV Canonical Form

Let $f_\alpha = \{ \psi_\alpha, \Gamma_\alpha \}$

where

$$\{ , \} = \frac{\partial}{\partial \underline{x}} \cdot \frac{\partial}{\partial \underline{p}} - \frac{\partial}{\partial \underline{p}} \cdot \frac{\partial}{\partial \underline{x}}$$

can show

$$[F, G] = \sum_x \int f_x \left(\frac{\partial}{\partial \underline{x}} \frac{\delta F}{\delta \underline{f}_x} \cdot \frac{\partial}{\partial \underline{p}} \frac{\delta G}{\delta \underline{f}_x} \right.$$

$$\left. - \frac{\partial}{\partial \underline{p}} \frac{\delta F}{\delta \underline{f}_x} \cdot \frac{\partial}{\partial \underline{x}} \frac{\delta G}{\delta \underline{f}_x} \right)$$

$$\rightarrow \sum_\alpha \int \left(\frac{\delta F}{\delta \psi_\alpha} \frac{\delta G}{\delta \Gamma_\alpha} - \frac{\delta G}{\delta \psi_\alpha} \frac{\delta F}{\delta \Gamma_\alpha} \right)$$

Use $\underline{E} \ \& \ \underline{A} \ (\nabla \times \underline{A} = \underline{B})$
for next.

